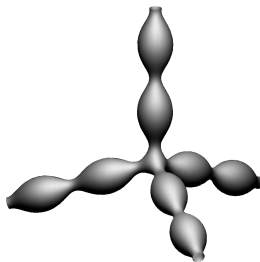


# CONSTANT MEAN CURVATURE $n$ -NOIDS WITH SYMMETRIES

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ABSTRACT. In this note I construct two families of constant mean curvature spheres with Delaunay ends via DPW. The construction uses symmetries of the DPW representation to unitarize the holonomies of the DPW potential.



## INTRODUCTION

In this note I construct two families of constant mean curvature spheres with Delaunay ends (noids) via DPW. The first family (“flowers”) consists of noids with  $n$  coplanar end axes and the symmetry of a  $n$ -gon slab. The members of the second family (“pods”) have  $n$  non-coplanar end axes and the symmetry of a  $(n - 1)$ -gon. The special case of the family of pods with four equal-weight ends have the symmetry of a tetrahedron.

The construction takes advantage of symmetries of the surface to unitarize the holonomies of the DPW potential. The symmetry of the potential induces a gauge symmetry via a loop  $g_0$  on the end monodromies  $M_j$ :

$$M_j = g_0^{-j} M_0 g_0^j.$$

$M_0$  and  $g_0$  are shown to be simultaneously unitarizable on  $\mathbb{S}^1$ , and so the  $M_j$  are simultaneously unitarizable. It follows that there exists a dressing which closes the ends of the surface.

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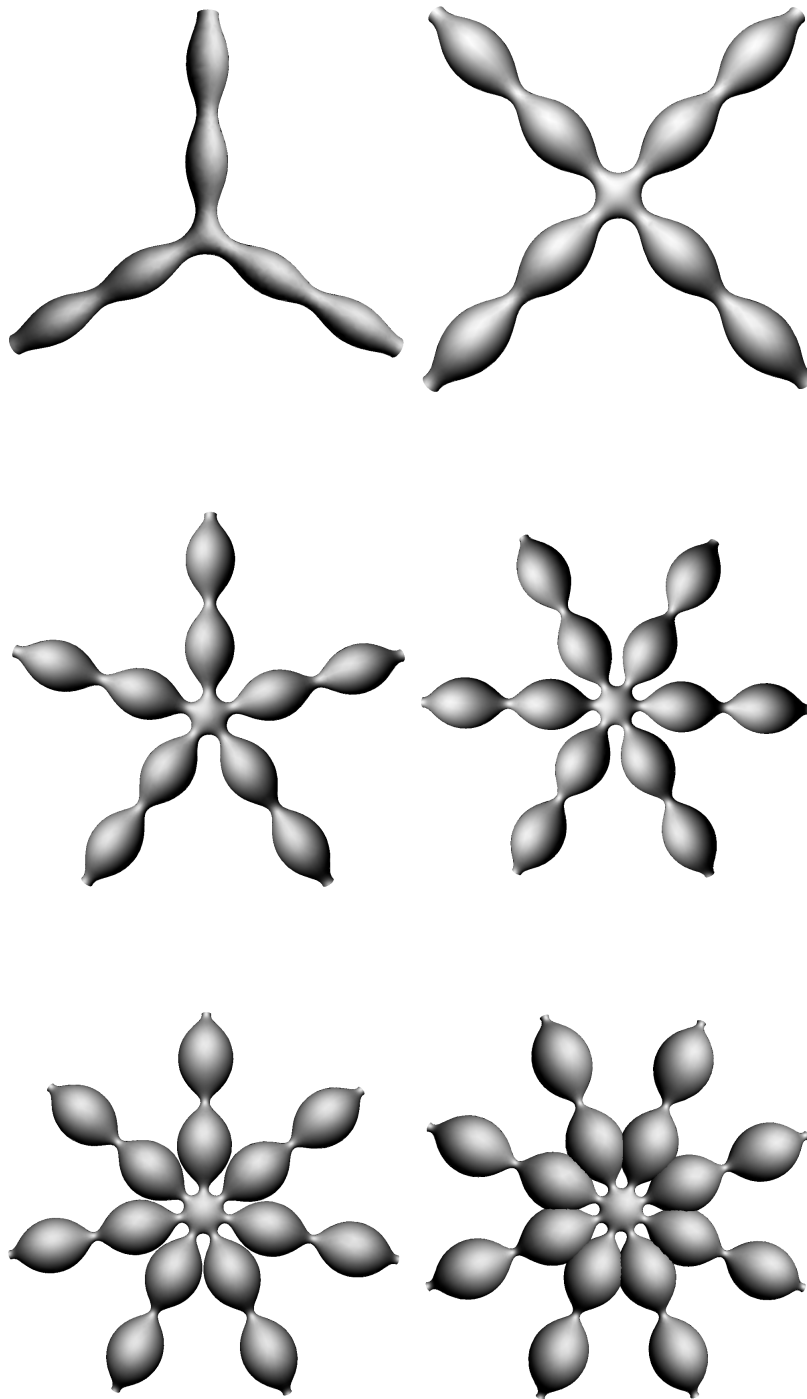


FIGURE 1. CMC  $n$ -flowers:  $n$ -noids with coplanar end axes, equal necksizes, and generically, a symmetry group of order  $4n$ .

## 1. SYMMETRIES

Lemmas 1.3 and 1.4 are basic results showing how a gauge symmetry of a potential descends to the ODE solution, its monodromy, the extended frame and the CMC immersion. The two lemmas detail the orientation-preserving and orientation-reversing cases respectively.

**Notation 1.1.** Let  $\Sigma$  be a Riemann surface. For a DPW potential  $\xi \in \Omega_1^1(\Lambda_\Sigma^{-1}\mathfrak{sl}_2(\mathbb{C}))$ , and a map  $g : \Sigma \rightarrow \Lambda_r\mathrm{SL}_2(\mathbb{C})$ , the *gauged potential* is

$$\xi \cdot g = g^{-1}dg + g^{-1}\xi g.$$

Gauging has the properties:

$$d\Phi = \Phi\xi \implies d(\Phi g) = (\Phi g)(\xi \cdot g)$$

and

$$(\xi \cdot g_1) \cdot g_2 = \xi \cdot (g_1 g_2).$$

**Notation 1.2.** Let  $\Sigma$  be a Riemann surface,  $\tilde{\Sigma}$  its universal cover, and  $\Gamma$  its group of deck transformations. Let  $z_0 \in \tilde{\Sigma}$ . Let  $\xi \in \Omega_1^1(\Lambda_\Sigma^{-1}\mathfrak{sl}_2(\mathbb{C}))$ . Let  $\Phi : \tilde{\Sigma} \rightarrow \Lambda_r\mathrm{SL}_2(\mathbb{C})$  be a solution to ODE  $d\Phi = \Phi\xi$ . The monodromy representation for  $\Phi$  relative to the basepoint  $z_0$  is the map  $M : \Gamma \rightarrow \Lambda_r\mathrm{SL}_2(\mathbb{C})$  defined by  $M_\sigma = \Phi(\sigma(z_0))\Phi(z_0)^{-1}$ .

**Lemma 1.3.** *Let  $\Sigma$  be a Riemann surface,  $\pi : \tilde{\Sigma} \rightarrow \Sigma$  its universal cover, and  $\Gamma$  its group of deck transformations. Let  $z_0 \in \tilde{\Sigma}$ . Let  $\xi \in \Omega_1^1(\Lambda_\Sigma^{-1}\mathfrak{sl}_2(\mathbb{C}))$  and let  $M : \Gamma \rightarrow \Lambda_r\mathrm{SL}_2(\mathbb{C})$  be the monodromy representation for  $\Phi$  relative to the basepoint  $z_0$ .*

*Let  $\tilde{\tau} : \tilde{\Sigma} \rightarrow \tilde{\Sigma}$  be a lift of an orientation-preserving automorphism  $\tau : \Sigma \rightarrow \Sigma$ . Let  $g : \Sigma \rightarrow \Lambda_r\mathrm{SL}_2(\mathbb{C})$  be a map such that*

$$(1.1) \quad \tau^*\xi = \xi \cdot g.$$

*and let  $V = ((\tilde{\tau}^*\Phi)(z_0))g^{-1}(z_0)\Phi^{-1}(z_0)$ .*

*Then (1)*

$$(1.2) \quad \tilde{\tau}^*\Phi = V\Phi g,$$

*and (2)*

$$(1.3) \quad \tilde{\tau}^*M = \mathrm{Ad}_V M.$$

*Assume  $V \in \Lambda_r^*\mathrm{SL}_2(\mathbb{C})$ . Then (3) with  $F = \mathrm{Uni}_r(\Phi)$ ,*

$$(1.4) \quad \tilde{\tau}^*F = VFD$$

*for some  $\lambda$ -independent diagonal matrix  $D$ , and (4) with  $f_\lambda = \mathrm{Sym}_\lambda[F]$ .*

$$(1.5) \quad \tilde{\tau}^*f_\lambda = \mathrm{Ad}_V f_\lambda - \frac{2}{H}V'V^{-1},$$

*so  $\tau^*f_\lambda$  is related to  $f_\lambda$  by an orientation-preserving isometry of  $\mathfrak{su}_2$ .*

*Proof.* The gauge symmetry (1.1) implies equation (1.2), since  $\tilde{\tau}^*\Phi$  and  $V\Phi g$  each satisfies the equation

$$\Psi^{-1}d\Psi = \tilde{\tau}^*\xi = \xi \cdot g$$

and they have the same value at  $z_0$ . Equation (1.3) follows.

Now assume  $V \in \Lambda_r^* \mathrm{SL}_2(\mathbb{C})$ . Let  $\Phi = FB$  be the Iwasawa factorization of  $\Phi$ . Then the Iwasawa factorization of  $\tilde{\tau}^*\Phi$  is  $(\tilde{\tau}^*F)(\tilde{\tau}^*B)$ , while that of  $V\Phi g$  is  $(VFD)(D^{-1}Bg)$  for some  $\lambda$ -independent diagonal matrix  $D$ . Equating unitary parts yields equation (1.4). The Sym formula yields equation (1.5).  $\square$

**Lemma 1.4.** *Let  $\Sigma$  be a Riemann surface,  $\pi : \tilde{\Sigma} \rightarrow \Sigma$  its universal cover, and  $\Gamma$  its group of deck transformations. Let  $z_0 \in \tilde{\Sigma}$ . Let  $\xi \in \Omega_1^1(\Lambda_{\tilde{\Sigma}}^{-1} \mathrm{sl}_2(\mathbb{C}))$  and let  $M : \Gamma \rightarrow \Lambda_r \mathrm{SL}_2(\mathbb{C})$  be the monodromy representation for  $\Phi$  relative to the basepoint  $z_0$ .*

*Let  $\iota$  be an orientation-reversing involution of the  $\lambda$ -plane. Let  $\tilde{\tau} : \tilde{\Sigma} \rightarrow \tilde{\Sigma}$  be a lift of an orientation-preserving automorphism  $\tau : \Sigma \rightarrow \Sigma$ . Let  $g : \Sigma \rightarrow \Lambda_r \mathrm{SL}_2(\mathbb{C})$  a map such that*

$$(1.6) \quad \overline{\tau^*\xi(\iota\lambda)} = \xi \cdot g.$$

*Let  $V(\lambda) = \overline{(\tilde{\tau}^*\Phi(\iota\lambda))(z_0)} g^{-1}(z_0, \lambda) \Phi^{-1}(z_0, \lambda)^{-1}$ .*

*Then (1)*

$$(1.7) \quad \overline{\tilde{\tau}^*\Phi(\iota\lambda)} = V\Phi g,$$

*and (2)*

$$(1.8) \quad \overline{\tilde{\tau}^*M(\iota\lambda)} = \mathrm{Ad}_V M(\lambda).$$

*Assume  $V \in \Lambda_r^* \mathrm{SL}_2(\mathbb{C})$  and  $\iota : \lambda \mapsto \bar{\lambda}$ . Then (3) with  $F = \mathrm{Uni}_r(\Phi)$ ,*

$$(1.9) \quad \overline{\tilde{\tau}^*F(\bar{\lambda})} = VFD$$

*for some  $\lambda$ -independent diagonal matrix  $D$ , and (4) with  $f_\lambda = \mathrm{Sym}_\lambda[F]$ .*

$$(1.10) \quad \overline{\tilde{\tau}^*f_\lambda} = -(\mathrm{Ad}_V f_\lambda - \frac{2}{H} V'V^{-1}).$$

*so  $\tilde{\tau}^*f_1$  is related to  $f_1$  by an orientation-reversing isometry of  $\mathrm{su}_2$ .*

*Proof.* The gauge symmetry (1.6) implies equation (1.7), because  $\overline{\tilde{\tau}^*\Phi(\iota\lambda)}$  and  $V\Phi g$  each satisfies the equation

$$\Psi^{-1}d\Psi = \overline{\tau^*\xi(\iota\lambda)} = \xi \cdot g$$

and they have the same value at  $z_0$ . Equation (1.8) follows.

Now assume  $V \in \Lambda_r^* \mathrm{SL}_2(\mathbb{C})$  and  $\iota : \lambda \mapsto \bar{\lambda}$ . Let  $\Phi = FB$  be the Iwasawa factorization of  $\Phi$ . Then the Iwasawa factorization of  $\overline{\tilde{\tau}^* \Phi(\bar{\lambda})}$  is

$$\overline{\tilde{\tau}^* \Phi(\bar{\lambda})} = (\overline{\tilde{\tau}^* F(\bar{\lambda})})(\overline{\tilde{\tau}^* B(\bar{\lambda})}),$$

while that of  $U\Phi g$  is  $(UFD)(D^{-1}Bg)$  for some  $\lambda$ -independent diagonal matrix  $D$ . Equating unitary parts yields equation (1.9). The Sym formula yields equation (1.10).  $\square$

**Remark 1.1.** A useful case of the above lemma is when  $\tilde{\tau} \circ \sigma = \sigma^{-1}$ , so  $\tilde{\tau}^* M_\sigma = M_{\tilde{\tau} \circ \sigma} = M_\sigma^{-1}$ .

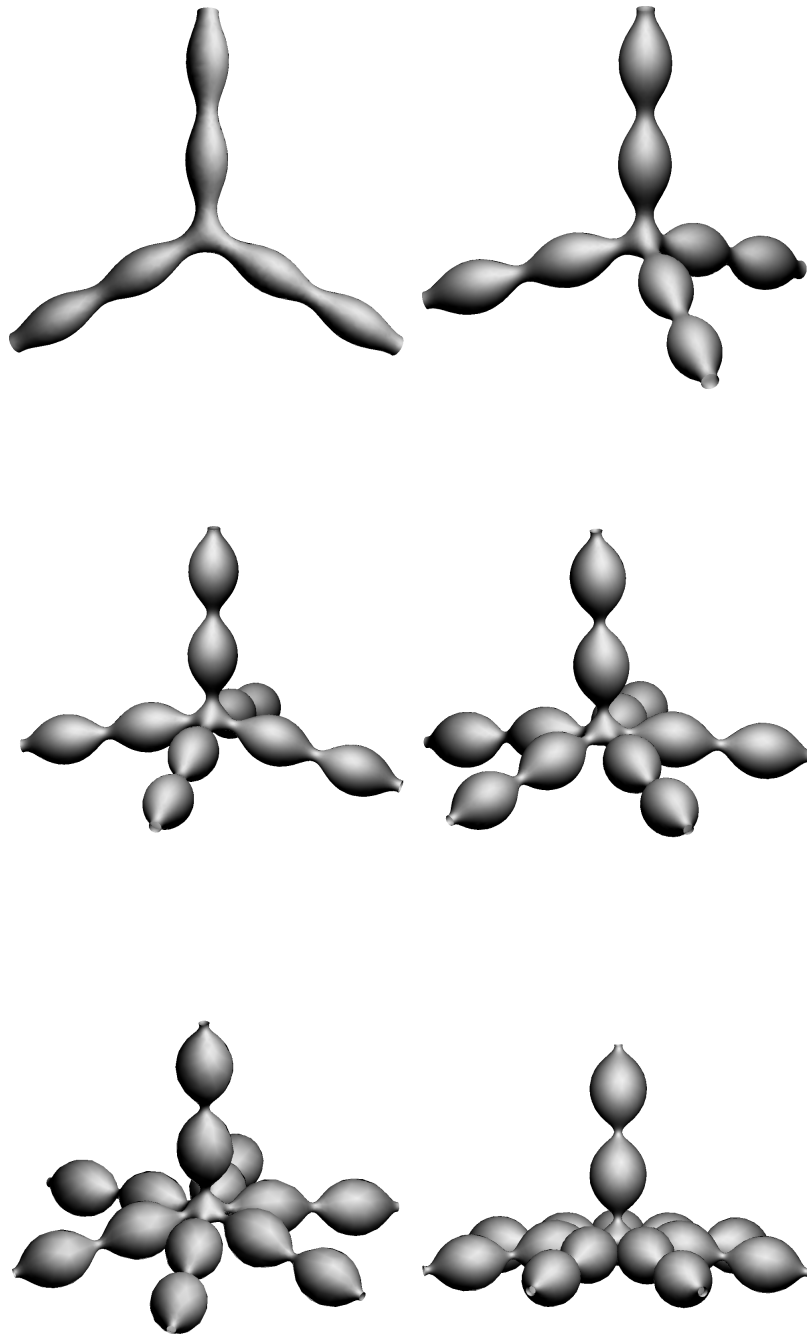


FIGURE 2. CMC  $n$ -pods:  $n$ -noids which generically have noncoplanar end axes, equal necksizes at  $n - 1$  ends, and a dihedral symmetry group of order  $2(n - 1)$ .

2.  $K$ -NOIDS

We now construct two families of  $n$ -noids. The first family contains noids with  $n$  equal-weight ends with coplanar axes. The noids in the second family have  $n - 1$  equal-weight ends, with generically noncoplanar end axes.

First we defined the DPW potential.

**Definition 2.1.** Let  $m \geq 1$  and let  $w_1 \in (-\infty, 1] \setminus \{0\}$ ,  $w_\infty \in (-\infty, 1]$  and  $v_k = \frac{1}{2}(1 - \sqrt{1 - w_k})$ ,  $k \in \{1, \infty\}$  satisfy

$$(2.1) \quad \begin{aligned} 0 &\leq m|\nu_1| + \nu_\infty \leq 1 \\ 0 &\leq m|\nu_1| - \nu_\infty \leq m - 1 \\ |w_\infty| &\leq m|w_1|. \end{aligned}$$

Let

$$(2.2) \quad \Sigma = \begin{cases} \mathbb{P}^1 \setminus (\{z^m = 1\} \cup \{\infty\}) & w_\infty \neq 0 \\ \mathbb{P}^1 \setminus \{z^m = 1\} & w_\infty = 0 \end{cases}$$

with standard conformal coordinate  $z$ . Define  $\xi_{m,w_1,w_\infty} \in \Omega_1^1(\Lambda_\Sigma^{-1}\mathfrak{sl}_2(\mathbb{C}))$  by

$$(2.3) \quad \xi_{m,w_1,w_\infty} = \begin{pmatrix} 0 & \lambda^{-1} \\ \frac{z^{m-2}(m^2w_1 + w_\infty(z^m - 1))}{16(z^m - 1)^2} & 0 \end{pmatrix} dz$$

**Theorem 2.2.** Let  $\xi_{n,w_1,w_\infty}$  as in definition 2.1. Then there exists a dressing matrix  $C$  such that the CMC immersion constructed by DPW from  $\xi_{m,w_1,w_\infty}$ , dressed by  $C$ , is a CMC noid.

In the case  $w_\infty \neq 0$ , the resulting CMC immersion has  $m + 1$  Delaunay ends:  $m$  weight  $w_1$  ends at the  $m$ 'th roots of unity and one weight  $w_\infty$  end at  $\infty$ . The surface has an order  $m - 2$  umbilic point at 0 and  $m$  order 1 umbilic points at the  $m$ 'th roots of  $1 - m^2w_1/w_\infty$ .

In the case  $w_\infty = 0$ , the resulting CMC immersion has  $m$  weight  $w_1$  Delaunay ends at the  $m$ 'th roots of unity. The potential  $\xi_{m,w_1,0}$  has an additional double pole at  $\infty$  which is a smooth point of the resulting CMC immersion. The surface has two order  $m - 1$  umbilics points at 0 and  $\infty$ .

*Proof.* The potential  $\xi = \xi_{m,w_1,w_\infty}$  has the following symmetry:

$$(2.4) \quad \rho : z \mapsto e^{2\pi i/m} z, \quad g = \begin{pmatrix} e^{-\pi i/m} & 0 \\ 0 & e^{\pi i/m} \end{pmatrix}, \quad \rho^* \xi = \xi \cdot g.$$

By lemma 1.3, the monodromies have the symmetry

$$M_j = g_0^{-j} M_0 g_0^j.$$

*Step 1: closing the ends.* We have that

$$(2.5) \quad \frac{1}{2} \operatorname{tr} g_0^{-1} = \cos \pi/m$$

$$(2.6) \quad \frac{1}{2} \operatorname{tr} M_0 = \cos 2\pi\nu_1$$

$$(2.7) \quad \frac{1}{2} \operatorname{tr} g^{-1} M_0 = \cos 2\pi(\frac{1}{2} - \nu_\infty)/m.$$

[NOTE: know third trace is root of unity. Fix which one by evaluation at  $\lambda = 1$ .] Hence  $g^{-1}$ ,  $M_0$  are simultaneously unitarizable on  $\mathbb{S}^1 \setminus \{\pm 1\}$  iff the spherical triangle inequalities hold for  $(1/(2m), \nu_1, (\frac{1}{2} - \nu_\infty)/m)$  on  $\mathbb{S}^1 \setminus \{\pm 1\}$ . A calculus argument shows that these are equivalent to inequalities (2.1).

By the gluing theorem [1], there exists an  $r$ -dressing  $C$  for  $r$  close to 1 for which the resulting immersion is closed at all its ends.  $\square$

**Corollary 2.3.** *Let  $\xi_{n,w_1,w_\infty}$  as in definition 2.1.*

*In the case  $w_\infty \neq 0$ , (1) for  $m > 2$  the end axes are non-coplanar; (2) (i) In the case that  $m > 3$ , or the case  $m = 3$  where not all necksizes are equal, the group of ambient isometries of the surface in  $\mathbb{R}^3$  is the dihedral group of a regular  $m$ -gon, of order  $2m$ ; (ii) In case of  $m = 3$  and equal necksizes, the ambient isometry group is that of the Platonic tetrahedron, of order 24.*

*In the case  $w_\infty = 0$ , (1) the end axes lie in a plane; (2) For  $m \geq 3$ , the group of ambient isometries of the surface in  $\mathbb{R}^3$  is that of a regular  $m$ -gon slab, with order  $4m$ .*

*Proof.* Using the symmetry of equation (2.4) and the fact that  $g_0^n = -I$ , lemma 1.3 implies that the surface has a orientation-preserving isometry induced by  $\rho$  of order  $n$ , which is hence a rotation.

The potential  $\xi = \xi_{m,w_1,w_\infty}$  has the following further symmetry:

$$(2.8) \quad \sigma : z \mapsto \bar{z}, \quad \overline{\sigma^* \xi(\bar{\lambda})} = \xi.$$

By lemma 1.4, the surface has an orientation-reversing isometry induced by  $\sigma$  of order 2 each, which is hence a reflection in a plane.

Now assume the special case  $n = 4$ ,  $w_1 = w_\infty$ , of four equal-weight ends. We show that the extrinsic symmetry group of the resulting surface is that of the regular tetrahedron. In this case,  $\xi = \xi_{4,w_1,w_1}$  has the additional symmetry:

$$(2.9) \quad v : z \mapsto \frac{z+2}{z-1}, \quad k = i\sqrt{3} \begin{pmatrix} z-1 & 0 \\ -\lambda & -3/(z-1) \end{pmatrix}, \quad \xi \cdot k = v^* \xi.$$

By lemma 1.3, the resulting CMC immersion has a order-2 rotation which exchanges two ends and fixes the other two.

Now take the case  $w_\infty = 0$ .  $\xi_{n,w_1,0}$  has the additional symmetry

$$(2.10) \quad \tau : z \mapsto 1/\bar{z}, \quad h = i \begin{pmatrix} z & 0 \\ -1 & -z^{-1} \end{pmatrix}, \quad \overline{\tau^* \xi(\bar{\lambda})} = \xi \cdot h.$$

By lemma 1.4, the surface has an orientation-reversing isometry induced by  $\tau$  of order 2, which is hence a reflection in a plane.  $\square$

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